

Classical Mechanics

- general equations for counterclockwise rotation

$$x(t) = R \cos(\omega t)$$

$$y(t) = R \sin(\omega t)$$

ω - angular frequency

$$\omega = \frac{2\pi}{T}$$

- Lagrangian: action of trajectory $A = \int_{t_0}^{t_1} (T - V) dt$

T - kinetic energy

V - potential energy

$$L = T - V \quad \text{Euler-Lagrange}$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{x}} - \frac{\partial L}{\partial x} = 0 \quad \text{or} \quad \frac{d}{dt} \frac{\partial L}{\partial \dot{x}} = \frac{\partial L}{\partial x}$$

$$p_i = \frac{\partial L}{\partial \dot{q}_i}$$

$$F_i = - \frac{\partial V}{\partial q_i}$$

$$L = L(q_i, \dot{q}_i, t)$$

p_i - canonical momentum

- Hamiltonian: if system is time translation invariant H conserved

$$H = T + V$$

$$\dot{p}_i = - \frac{\partial H}{\partial q_i}$$

$$\dot{q}_i = \frac{\partial H}{\partial p_i}$$

$$H = H(q_i, p_i)$$

Gibbs-Liouville theorem - divergence of phase space of H is 0

$$\vec{\nabla} \cdot \vec{V} = \sum_i \left(\frac{\partial}{\partial q_i} \frac{\partial H}{\partial p_i} - \frac{\partial}{\partial p_i} \frac{\partial H}{\partial q_i} \right) = 0$$

- phase space of H $2N$ dimensional (p, q for every N)

- Poisson brackets - $\dot{F} = \{F, H\} = \sum_i \left(\frac{\partial F}{\partial q_i} \frac{\partial H}{\partial p_i} - \frac{\partial F}{\partial p_i} \frac{\partial H}{\partial q_i} \right)$ also $\dot{q}_i = \{q_i, H\}$

Definition + arguments:

$$\dot{p}_i = \{p_i, H\}$$

$$\{A, C\} = \sum_i \left(\frac{\partial A}{\partial q_i} \frac{\partial C}{\partial p_i} - \frac{\partial A}{\partial p_i} \frac{\partial C}{\partial q_i} \right)$$

$$\{A, C\} = -\{C, A\} \quad \{A, A\} = 0 \quad \{kA, C\} = k\{A, C\}$$

$$\{(A+B), C\} = \{A, C\} + \{B, C\}$$

$$\{AB, C\} = B\{A, C\} + A\{B, C\}$$

$$\text{also: } \{q_i, q_j\} = 0$$

$$\{p_i, p_j\} = 0$$

$$\{q_i, p_j\} = \delta_{ij}$$

$$\{q^n, p\} = n q^{(n-1)}$$

δ_{ij} - Kronecker delta function $\delta_{ij} = 0$ if $i \neq j$ $\delta_{ij} = 1$ if $i = j$

Poisson bracket of quantity with another one which is invariant in some translation gives you the derivative of first quantity with respect to the property in which second shows translational invariance eg.) H - time translational invariance

if G conserved in time translations

$$\{G, H\} = 0$$

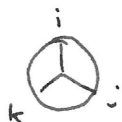
$$\frac{d}{dt} f = \dot{f} = \{f, H\}$$

$$\text{or } \{x, L_z\} \propto \delta x$$

$$\{y, L_z\} \propto \delta y$$

$$\{z, L_z\} \propto \delta z$$

Levi-Civita symbol ϵ_{ijk}



if $i=j=k$ $\epsilon_{ijk} = 0$

if order clockwise ϵ_{ijk} (i,j,k / j,k,i / k,i,j)
 $\epsilon_{ijk} = 1$

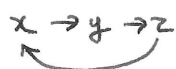
if order anti clockwise ϵ_{ijk} (i,k,j / k,j,i / j,i,k)
 $\epsilon_{ijk} = -1$

Del ∇ $\nabla_x \equiv \frac{\partial}{\partial x}$ $\nabla_y \equiv \frac{\partial}{\partial y}$ $\nabla_z \equiv \frac{\partial}{\partial z}$

↳ gradient of scalar A is a vector $\vec{\nabla} A = \frac{\partial A}{\partial x}, \frac{\partial A}{\partial y}, \frac{\partial A}{\partial z}$

↳ divergence of vector $\vec{A}$ is a scalar $\vec{\nabla} \cdot \vec{A} = \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z}$

↳ curl of vector $\vec{A}$ is a vector $\vec{\nabla} \times \vec{A}$



$$(\vec{\nabla} \times \vec{A})_x = \frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z}$$

$$(\vec{\nabla} \times \vec{A})_y = \frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x}$$

$$(\vec{\nabla} \times \vec{A})_z = \frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y}$$

curl has no divergence so $\vec{\nabla} \cdot (\vec{\nabla} \times \vec{A}) = 0$

- magnetic field - 0 divergence so represented by vector potential

$$\vec{B} = \vec{\nabla} \times \vec{A} \quad B_x = \frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z}$$

By ...

gauge transformation changes $\vec{A}$ but $\vec{B}$ not affected so $\vec{B}$ gauge transformation invariant

transformation - gradient function of scalar

$$\vec{A}' = \vec{A} + \vec{\nabla} \alpha$$

electric field $\vec{E} = -e \vec{\nabla} V$

magnetic field (Lorentz force)

$$\vec{F} = \frac{e}{c} \vec{v} \times \vec{B}$$

- gravity

$$\vec{F}_{\text{grav}} = F_{\text{gm}} \hat{r}$$

$$F_{\text{grav}} = - \frac{G m_1 m_2}{r^2}$$

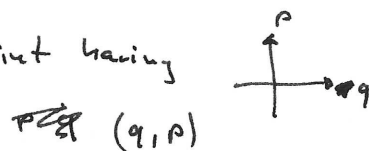
$$\hat{r} = \frac{\vec{r}}{|\vec{r}|}$$

$$= - \frac{G m_1 m_2}{r^2} \hat{r}$$

$r = |\vec{r}| \leftarrow$ or just r (length of radius)

- phase space - state space of mechanical system with each point having unique position value and momentum value

- 2N dimensional



- configuration space - in classical mechanics it is the vector space defined by generalised coordinates describing state of system

- euclidean or polar

- N dimensional

- state space - set of all possible configurations (states) of a system

- variables on axis are state variables

- Hessian matrix

$$H = \begin{pmatrix} \frac{\partial^2 A}{\partial x^2} & \frac{\partial^2 A}{\partial x \partial y} \\ \frac{\partial^2 A}{\partial y \partial x} & \frac{\partial^2 A}{\partial y^2} \end{pmatrix}$$

$$\frac{\partial^2 A}{\partial x \partial y} = \frac{\partial^2 A}{\partial y \partial x}$$

$$\text{Det } H = \frac{\partial^2 A}{\partial x^2} \frac{\partial^2 A}{\partial y^2} - \frac{\partial^2 A}{\partial y \partial x} \frac{\partial^2 A}{\partial x \partial y} = \frac{\partial^2 A}{\partial x^2} \frac{\partial^2 A}{\partial y^2} - 2 \frac{\partial^2 A}{\partial x \partial y}$$

$$\text{Tr } H = \frac{\partial^2 A}{\partial x^2} + \frac{\partial^2 A}{\partial y^2}$$

local minimum if : $\text{Det } H > 0$ $\text{Tr } H > 0$ $\therefore \text{min}$

local maximum if : $\text{Det } H > 0$ $\text{Tr } H < 0$ $\therefore \text{max}$

saddle point if : $\text{Det } H < 0$ $\therefore \text{point of inflexion}$